

SR as a meta-theory: **laws of physics must be Lorentz-invariant.**

Laws of classical mechanics are not: they are invariant under Galilean transformations, not Lorentz transformations.

=> we need a revised, relativistic, mechanics!

Starting point: Assume each particle has an invariant **proper mass**
(or rest mass) m_0

Define **4-momentum** $\mathbf{P} = m_0 \mathbf{U}$

Basic axiom: conservation of \mathbf{P} in particle collisions: $\Sigma^* \mathbf{P}_n = 0$

(in the sum, \mathbf{P} of each particle going into the collision is counted positively while \mathbf{P} of each particle coming out is counted negatively)

Since a sum of 4-vectors is a 4-vector, this equation is Lorentz-invariant.

$$\mathbf{P} = m_0 \mathbf{U} = m_0 \gamma(u) (\mathbf{u}, c) \equiv (\mathbf{p}, mc)$$

relativistic mass: $m = \gamma(u) m_0$

relativistic momentum: $\mathbf{p} = m \mathbf{u}$

Conservation of 4-momentum yields both: $\Sigma^* \mathbf{p} = 0$ $\Sigma^* m = 0$

These are Newtonian conservation laws when $c \rightarrow \infty$

What about conservation of energy?

When $v/c \ll 1$: $m = m_0(1 - v^2/c^2)^{-1/2} \approx m_0 \left(1 + \frac{1}{2} \frac{v^2}{c^2} \right)$

$$mc^2 \approx m_0c^2 + \frac{1}{2}m_0v^2$$

Einstein proposed the equivalence of mass and energy: $E = mc^2$

$$\Rightarrow \mathbf{P} = (\mathbf{p}, E/c)$$

Relativistic energies

rest energy: $m_0 c^2$

kinetic energy: $T = mc^2 - m_0 c^2 = (\gamma - 1) m_0 c^2$

Note: Potential energy associated with position of particle in an external electromagnetic or gravitational field does not contribute to the relativistic mass of a particle. When particle's PE changes, it's actually the energy of the field that's changing.

H atom: energy of electron = -13.6 eV (binding energy)

$$\Rightarrow m_{\text{atom}} c^2 = m_p c^2 + m_e c^2 - 13.6 \text{ eV}$$

\Rightarrow atom has less mass than its constituents!

$$1 \text{ eV} = 1.602 \times 10^{-12} \text{ ergs} = 1.602 \times 10^{-19} \text{ J}$$

$$m_p = 1.67 \times 10^{-24} \text{ g} \quad ; \quad m_e = 9.11 \times 10^{-28} \text{ g}$$

$$\Rightarrow \Delta m / m = 13.6 \text{ eV} / (m_p + m_e) c^2 \approx 1.8 \times 10^{-8}$$

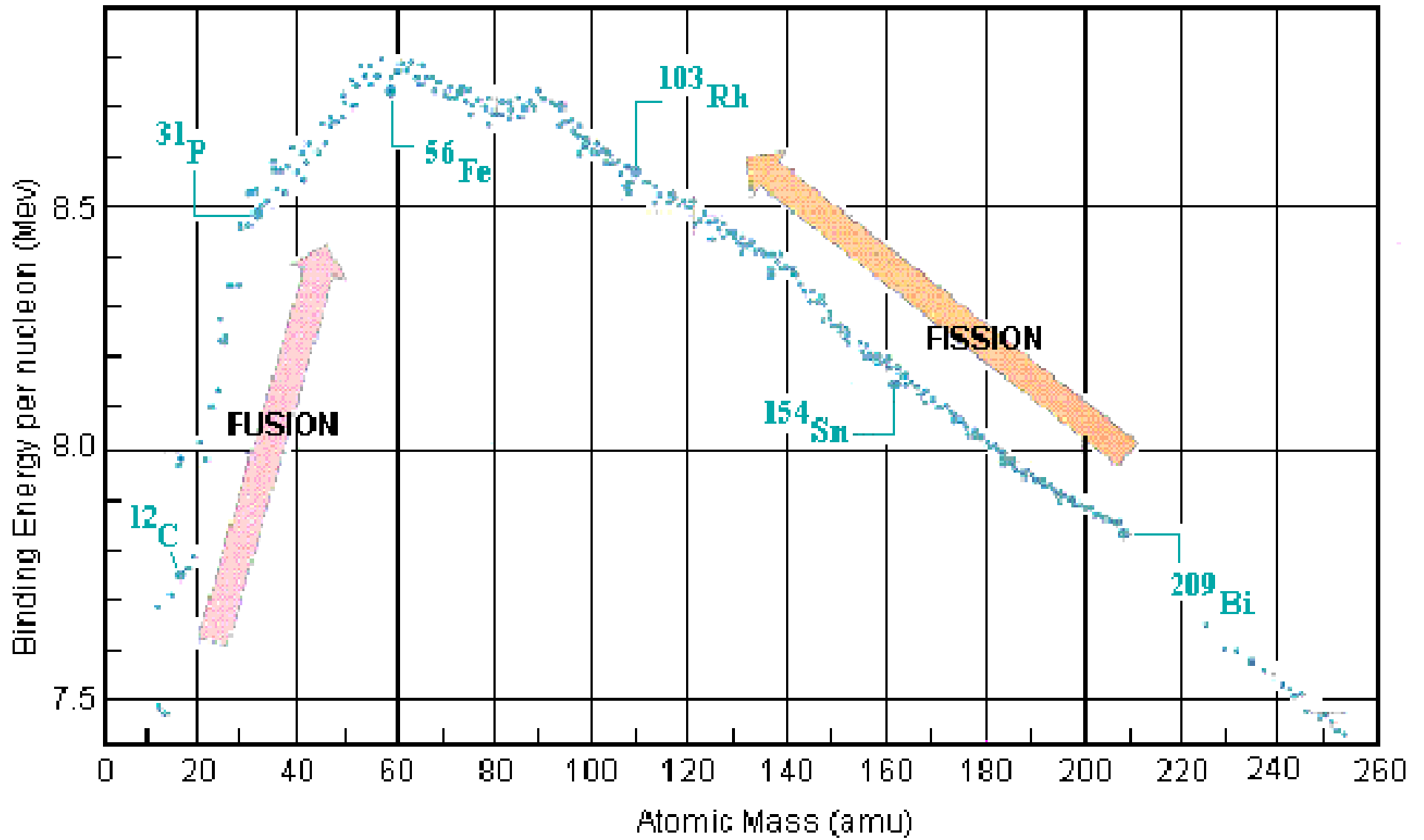
Nuclear physicists use the **atomic mass unit** u, defined so that the mass of the ^{12}C atom is exactly 12 u.

$$1 \text{ u} = 931.502 \text{ MeV} / c^2$$

$$m_p = 1.00727647 \text{ u} \quad ; \quad m_n = 1.00866501 \text{ u} \quad ; \quad m_e = 5.485803 \times 10^{-4} \text{ u}$$

$$\Rightarrow \Delta m / m \approx 0.8\% \quad \text{for the } ^{12}\text{C} \text{ nucleus}$$

$$\text{Binding energy} \approx 90 \text{ MeV} \quad (\approx 7.5 \text{ MeV per nucleon})$$



$$\mathbf{P} = m_0 \gamma(u) (\mathbf{u}, c) = (\mathbf{p}, mc) = (\mathbf{p}, E/c)$$

$$\Rightarrow \mathbf{P}^2 = E^2/c^2 - \mathbf{p}^2$$

In particle's rest frame: $\mathbf{P} = (\mathbf{0}, m_0 c) \Rightarrow \mathbf{P}^2 = m_0^2 c^2$

$$\Rightarrow E^2/c^2 - \mathbf{p}^2 = m_0^2 c^2 \Rightarrow E^2 = m_0^2 c^4 + \mathbf{p}^2 c^2$$

For a photon: $m_0 = 0 \Rightarrow E = pc$, $\mathbf{P} = E/c (\mathbf{n}, 1)$

$$E = h\nu = hc/\lambda \quad , \quad \mathbf{n} = \text{a unit vector in dir of propagation}$$

Since \mathbf{P} is a 4-vector, E and \mathbf{p} transform as follows under a standard LT:

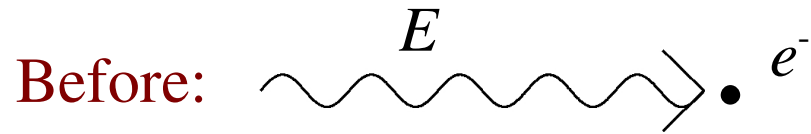
$$p_x' = \gamma(p_x - \beta E/c) \quad E' = \gamma(E - \beta p_x c)$$

$$p_y' = p_y$$

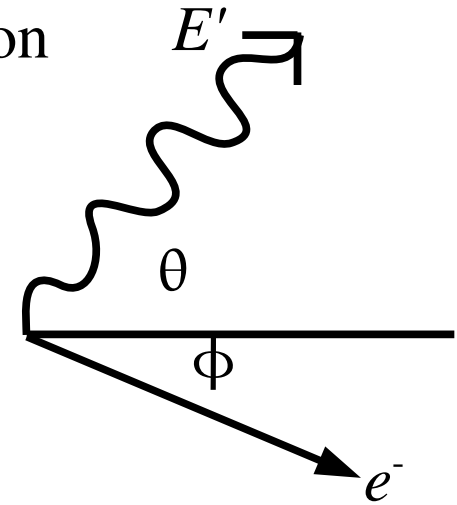
$$p_z' = p_z$$

Applications

1. Compton scattering: photon strikes stationary electron



After:



$$\mathbf{P}_{e,i} = m_e c (\mathbf{0}, 1) \quad \mathbf{P}_{\gamma,f} = \frac{E'}{c} (\hat{x} \cos \theta + \hat{y} \sin \theta, 1)$$

$$\mathbf{P}_{\gamma,i} = \frac{E}{c} (\hat{x}, 1)$$

$$\mathbf{P}_{e,i} + \mathbf{P}_{\gamma,i} = \mathbf{P}_{e,f} + \mathbf{P}_{\gamma,f}$$

$$(\mathbf{P}_{e,i} + \mathbf{P}_{\gamma,i} - \mathbf{P}_{\gamma,f})^2 = \mathbf{P}_{e,f}^2 = m_e^2 c^2$$

$$m_e^2 c^2 + (\mathbf{P}_{\gamma,i} - \mathbf{P}_{\gamma,f})^2 + 2\mathbf{P}_{e,i} \cdot (\mathbf{P}_{\gamma,i} - \mathbf{P}_{\gamma,f}) = m_e^2 c^2$$

$$-2\mathbf{P}_{\gamma,i} \cdot \mathbf{P}_{\gamma,f} + 2\mathbf{P}_{e,i} \cdot \mathbf{P}_{\gamma,i} - 2\mathbf{P}_{e,i} \cdot \mathbf{P}_{\gamma,f} = 0$$

$$\mathbf{P}_{e,i} = m_e c (\mathbf{0}, 1)$$

$$\mathbf{P}_{\gamma,i} = \frac{E}{c} (\hat{x}, 1)$$

$$\mathbf{P}_{\gamma,f} = \frac{E'}{c} (\hat{x} \cos \theta + \hat{y} \sin \theta, 1)$$

$$\mathbf{P}_{e,i} \cdot \mathbf{P}_{\gamma,i} = \mathbf{P}_{e,i} \cdot \mathbf{P}_{\gamma,f} + \mathbf{P}_{\gamma,i} \cdot \mathbf{P}_{\gamma,f}$$

$$m_e E = m_e E' + \frac{EE'}{c^2} (1 - \cos \theta)$$

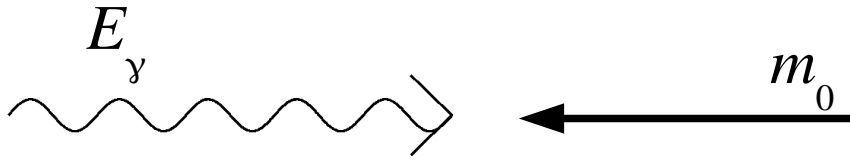
$$\frac{hc}{\lambda} = \frac{hc}{\lambda'} + \frac{h^2 c^2 (1 - \cos \theta)}{m_e c^2 \lambda \lambda'}$$

$$\lambda' = \lambda + \frac{h}{m_e c} (1 - \cos \theta)$$

$h / m_e c$ is called the
 “Compton wavelength”
 of the electron.

2. Inverse Compton scattering: photon scatters off a highly energetic charged particle

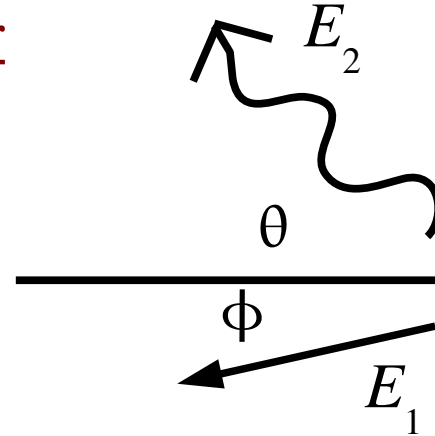
Before



$$\mathbf{P}_{p,i} = (-p\hat{x}, E/c) \quad \mathbf{P}_{\gamma,i} = \frac{E_{\gamma}}{c} (\hat{x}, 1)$$

$$\mathbf{P}_{\gamma,f} = \frac{E_2}{c} (-\hat{x} \cos \theta + \hat{y} \sin \theta, 1)$$

After



As before: $\mathbf{P}_{p,i} \cdot \mathbf{P}_{\gamma,i} = \mathbf{P}_{p,i} \cdot \mathbf{P}_{\gamma,f} + \mathbf{P}_{\gamma,i} \cdot \mathbf{P}_{\gamma,f}$

$$\Rightarrow E_{\gamma}(E + pc) = E_2(E - pc \cos \theta) + E_{\gamma}E_2(1 + \cos \theta)$$

Assume $E \gg m_0 c^2 \gg E_{\gamma}$

$$pc = (E^2 - m_0^2 c^4)^{1/2} \approx E \left(1 - \frac{m_0^2 c^4}{2E^2} \right) \quad (E \gg m_0 c^2)$$

$$E_\gamma(E + pc) = E_2(E - pc \cos \theta) + E_\gamma E_2(1 + \cos \theta) \quad \text{(from previous slide)}$$

$$\begin{aligned} 2EE_\gamma &\approx E_2 \left(E - E \cos \theta + \frac{m_0^2 c^4}{2E} \cos \theta + E_\gamma + E_\gamma \cos \theta \right) \\ &= E_2 \left[(E + E_\gamma) - \left(E - E_\gamma - \frac{m_0^2 c^4}{2E} \right) \cos \theta \right] \end{aligned}$$

Maximum energy is transferred to photon when $\cos \theta = 1$:

$$2EE_\gamma \approx 2E_2E_\gamma \left(1 + \frac{m_0^2 c^4}{4EE_\gamma} \right) \quad \Rightarrow \quad E_2 = \frac{E}{1 + m_0^2 c^4 / 4EE_\gamma}$$

Cosmic Microwave Background (CMB): 3 K blackbody spectrum

=> typical photon energy $\sim kT = 3 \times 10^{-4}$ eV

Very high energy cosmic ray proton: 10^{20} eV ; $m_p c^2 = 938$ MeV

=> CMB photon can be upscattered to $\sim 10^{19}$ eV (source of gamma rays)

3. Particle accelerators: “available energy” for creating new particles is the energy in the CM frame

Consider 2 arrangements:

a) a 30 GeV proton collides with a proton at rest

b) two 15 GeV protons collide head-on

(lab and CM frames are identical => 30 GeV of available energy)

a) lab frame: $\mathbf{P}_i = [\gamma\beta m_p c, (\gamma+1) m_p c]$ CM frame: $\mathbf{P} = (\mathbf{0}, E/c)$

$$\mathbf{P}^2 = (\gamma + 1)^2 m_p^2 c^2 - \gamma^2 \beta^2 m_p^2 c^2 = E^2 / c^2$$

$$[(\gamma^2 - \gamma^2 \beta^2) + 1 + 2\gamma] m_p^2 c^2 = E^2 / c^2$$

$$\begin{aligned} E^2 &= 2(1 + \gamma) m_p^2 c^4 \\ &= 2m_p c^2 [m_p c^2 + \gamma m_p c^2] \\ &= 2(0.938 \text{ GeV}) [0.938 \text{ GeV} + 30 \text{ GeV}] \end{aligned}$$

$$\Rightarrow E = 7.6 \text{ GeV}$$

To reach an available energy of 30 GeV in arrangement (a), we would need a 480 GeV proton!

4. Cosmic ray cutoff: high-energy nucleons can undergo the reaction $\gamma + N \rightarrow N + \pi$

Reactions with CMB photons ($E \sim 3 \times 10^{-4}$ eV) probably produces a cutoff in the cosmic ray spectrum: **no CRs with energy above the threshold energy for the reaction.**

Assuming a head-on collision, what's the threshold energy?

We need $E = (m_N + m_\pi) c^2$ in CM frame $\Rightarrow \mathbf{P}_f = (\mathbf{0}, m_N c + m_\pi c)$

Lab frame, before collision: $\mathbf{P}_\gamma = \frac{E_\gamma}{c}(\hat{x}, 1)$ $\mathbf{P}_N = (-p_N \hat{x}, E_N/c)$

$$\Rightarrow \mathbf{P}_i = \left(\frac{E_\gamma}{c} - p_N, 0, 0, \frac{E_\gamma + E_N}{c} \right)$$

$$(E_\gamma + E_N)^2 - (E_\gamma - p_N c)^2 = (m_N + m_\pi)^2 c^4$$

$$(E_\gamma + E_N)^2 - (E_\gamma - p_N c)^2 = (m_N + m_\pi)^2 c^4$$

$$E_N^2 + 2E_\gamma E_N - p_N^2 c^2 + 2E_\gamma p_N c = (m_N + m_\pi)^2 c^4$$

$$2E_\gamma(E_N + p_N c) = [(m_N + m_\pi)^2 - m_N^2] c^4$$

$$E_N + p_N c = \frac{m_\pi(2m_N + m_\pi)c^4}{2E_\gamma} \approx 6 \times 10^{14} \text{ MeV}$$

$$(m_N c^2 \approx 940 \text{ MeV}, m_\pi c^2 \approx 140 \text{ MeV})$$

$$\Rightarrow E_N \gg m_N c^2 \Rightarrow \text{LHS} \rightarrow 2 E_N \Rightarrow E_N \approx 3 \times 10^{14} \text{ MeV}$$

What about forces?

Define **4-force** $\mathbf{F} = d\mathbf{P}/d\tau = d(m_0\mathbf{U}) / d\tau = m_0\mathbf{A} + (dm_0/d\tau)\mathbf{U}$

Limited version of Newton's 3rd Law for contact collisions between 2 particles: **during contact, τ is the same for both particles.**

$$\mathbf{F}_1 + \mathbf{F}_2 = d(\mathbf{P}_1 + \mathbf{P}_2) / d\tau = 0 \quad \Rightarrow \quad \mathbf{F}_2 = -\mathbf{F}_1$$

Relativistic 3-force $\mathbf{f} = d\mathbf{p}/dt = d(m\mathbf{u})/dt$

$$\mathbf{F} = \frac{d\mathbf{P}}{d\tau} = \gamma(u) \frac{d}{dt} (\mathbf{p}, E/c) = \gamma(u) \left(\mathbf{f}, \frac{1}{c} \frac{dE}{dt} \right)$$

\mathbf{f} becomes Newtonian force when $\beta \rightarrow 0$

4th component of \mathbf{F} is proportional to the power absorbed by the particle.

$$\mathbf{F} \cdot \mathbf{U} = \left(m_0 \mathbf{A} + \frac{dm_0}{d\tau} \mathbf{U} \right) \cdot \mathbf{U} = c^2 \frac{dm_0}{d\tau}$$

$\Rightarrow \mathbf{F} \cdot \mathbf{U}$ = proper rate at which particle's internal energy increases

$$\begin{aligned} \text{Also: } \mathbf{F} \cdot \mathbf{U} &= \gamma(u) \left(\mathbf{f}, \frac{1}{c} \frac{dE}{dt} \right) \cdot \gamma(u) (\mathbf{u}, c) \\ &= \gamma^2(u) \left(\frac{dE}{dt} - \mathbf{f} \cdot \mathbf{u} \right) \end{aligned}$$

Equating above 2 expressions for $\mathbf{F} \cdot \mathbf{U} \Rightarrow$ **rest mass-preserving forces are characterized by:**

$$\mathbf{F} \cdot \mathbf{U} = 0$$

$$\mathbf{f} \cdot \mathbf{u} = dE/dt \quad \Rightarrow \quad \mathbf{f} \cdot d\mathbf{r} = dE$$

$$\mathbf{F} = \gamma(u) (\mathbf{f}, \mathbf{f} \cdot \mathbf{u} / c)$$

(as in Newtonian theory; the added energy is entirely kinetic)

4-force transformation: $(Q \equiv dE/dt)$

$$f'_1 = \frac{f_1 - vQ/c^2}{1 - u_1v/c^2} \left(= \frac{f_1 - v\mathbf{f} \cdot \mathbf{u}/c^2}{1 - u_1v/c^2}, \quad m_0 = \text{const} \right)$$

$$f'_2 = \frac{f_2}{\gamma(v)(1 - u_1v/c^2)}$$

$$f'_3 = \frac{f_3}{\gamma(v)(1 - u_1v/c^2)}$$

$$Q' = \frac{Q - vf_1}{1 - u_1v/c^2}$$

Note: for rest mass-preserving force, \mathbf{f} is invariant among IFs with relative velocities parallel to force.

For a rest mass-preserving force:

$$\begin{aligned}
 \mathbf{f} &= \frac{d(m\mathbf{u})}{dt} = m\mathbf{a} + \frac{dm}{dt}\mathbf{u} \\
 &= \gamma m_0 \mathbf{a} + \frac{1}{c^2} \frac{dE}{dt} \mathbf{u} \\
 &= \gamma m_0 \mathbf{a} + \frac{\mathbf{f} \cdot \mathbf{u}}{c^2} \mathbf{u}
 \end{aligned}$$

\mathbf{a} is not generally
parallel to \mathbf{f}

$$\mathbf{f} \parallel \mathbf{u} : \quad f_{\parallel} = \gamma m_0 a_{\parallel} + f_{\parallel} u^2 / c^2$$

$$\Rightarrow f_{\parallel} = \gamma^3 m_0 a_{\parallel}$$

$$\mathbf{f} \perp \mathbf{u} : \quad f_{\perp} = \gamma m_0 a_{\perp}$$

A moving particle offers
different inertial resistances
to the same force, depending
on whether it is applied
longitudinally or transverse.

Also for a rest mass-preserving force:

$$\mathbf{F} = m_0 \mathbf{A} + dm_0/d\tau \mathbf{U} = m_0 d^2 \mathbf{R}/d\tau^2 = \gamma(u) (\mathbf{f}, c^{-1} dE/dt)$$

$$\Rightarrow \gamma(u) \mathbf{f} = m_0 d^2 \mathbf{r}/d\tau^2$$